

Совещание по использованию рассеяния нейтронов и синхротронного излучения в конденсированных средах



Эффект нейтронного гетеродинирования и быстрая кинетика нано-систем

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Where the spins are and what they are doing in nano-materials. Current subjects for reflectometry:

(c)

(d)

(e)

F - pinned

 $2\Lambda_{\rm chem}$

 Λ_{chem}

 Λ_{mag}

a) Exchange coupled
bi-layers and superlattice with
antiferromagnetic ordering
(GMR and TMR systems)

(a)

(b)

Ga_{1-x}Mn_xAs

GaAs

c) Laterally patterned magnetic films, nano-wires, etc.

d) Ferromagnetic films
on antiferromagnetic
substrates with
Exchange bias
through common interfaces

b) Dilute magnetic semiconductors as spin-injectors in semiconductor heterostructures e) Spring magnets: soft magnetic layer exchange coupled to a magnetically hard layer (spin valves)

Also: Superconducting films, polymer films, membranes, etc

Recent Review Articles

H. Zabel, K. Theis-Bröhl, B.P. Toperverg,

"Polarized neutron reflectivity and scattering of magnetic nanostructures and spintronic materials"

Handbook of Magnetism and Advanced Magnetic Materials, H. Kronmüller & S. Parkin (Eds.), NY, Wiley 2007, pp. 1237-1288

<u>A. Remhof, A. Westphalen, K. Theis-Bröhl, J. Grabis, A. Nefedov,</u> B. Toperverg, H. Zabel

"Magnetization Reversal Studies of Periodic Magnetic Arrays via Scattering Methods"

Springer Series in materials science 94, Springer-Verlag Berlin Heidelberg 2007 pp.65 - 97

H.-J.C. Lauter, V. Lauter, B.P. Toperverg

"Reflectivity, Off-Specular Scattering, and GI-SAS: Neutrons"

Polymer Science: A Comprehensive Reference, Vol 2, pp. 411–432 (2012)

Matyjaszewski K & Möller M (eds.) Amsterdam: Elsevier BV

M.R. Fizsimmons & I.K. Schuller

"Neutron scattering – The key characterization tool for nanostructured magneic materials" JMMM, 350 (2014) 199-208

To-day spintronic material application: Magnetic Random Access Memory (MRAM)

- \rightarrow "read-write" 2D arrays of spin valves
 - Writing with weak magnetic field
- Reading with electric current

MRAM-Storage Device Soft magnetic layer Connecting - metal bars **Tunnel** barrier Hard magnetic layer **[0]** Contacting metal bars

Neutrons: Intrinsic mechanisms and rates of spin rearrangement in nano-elements and their ensembles Emerging Challenges: non-linear spin dynamics & kinetics

Magnetic Domain-Wall Racetrack Memory

Stuart S. P. Parkin,* Masamitsu Hayashi, Luc Thomas





Magnetic ratchet for <u>three-dimensional spintronic</u> memory and logic

Reinoud Lavrijsen¹, Ji-Hyun Lee¹, Amalio Fernández-Pacheco¹, Dorothée C. M. C. Petit¹, Rhodri Mansell¹ & Russell P. Cowburn¹

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Layer-by-layer time resolved magnetometry with PNR?

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Frequency dependence of magnetization reversal in thin Fe(100) films

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Reduced magnetization $m(t) = M(t) / M_{sat} = w_{\uparrow} - w_{\downarrow} = 1 - 2w_{\downarrow}(t)$,

Surface fraction $w_{\downarrow}(t) = \frac{1}{l} \int_{0}^{t} dt' v(t')$ with negative magnization, *l* is domain size

 $v(t) = \mu_{\rm DW} H(t) \text{ is DW velocity (Landau & Lifshits, 1935), } \mu_{\rm DW} \text{ is DW mobility}$ If $H(t) = H_0 \sin(\omega \cdot t)$ then $w_{\downarrow}(t) = \mu_{\rm DW} H_0 \frac{1 - \cos(\omega \cdot t)}{\omega \cdot l}$, $m(t) = \overline{m} + \Delta m(t)$, where $\overline{m} = 1 - 2\mu_{\rm DW} H_0 \frac{1}{\omega \cdot l}$, $\Delta m(t) = 2\mu_{\rm DW} H_0 \frac{\cos(\omega \cdot t)}{\omega \cdot l}$

Landau-Lifshits regime vs Walker break down



Schyer & Walker, 1974

Reflection for polarization collinear with magnetization



Reflection coefficients (ideal polarization, no analysis)

$$\mathbf{R}(\vec{b}) = \frac{1}{2} = \left| \{\mathbf{R}_{+} \mid^{2} \cdot \left[+ (\vec{P} \cdot \vec{b}) + |\mathbf{R}_{-}|^{2} \cdot \left[- (\vec{P} \cdot \vec{b}) \right] \right] \right|$$

 R_{\pm} are reflection amplitudes for \pm spin states

Surface average for time moment *t* :

$$\mathbf{R} = w_{\uparrow} \mathbf{R}(\vec{b}) + w_{\downarrow} \mathbf{R}(-\vec{b}), \qquad \vec{b} = \vec{B}_{\parallel} / |\vec{B}_{\parallel}|$$
$$\mathbf{R}^{+}(t) = \frac{1}{2} |\vec{R}_{+}|^{2} \cdot [+m(t)] + |\vec{R}_{-}|^{2} \cdot [-m(t)],$$
$$\mathbf{R}^{-}(t) = \frac{1}{2} |\vec{R}_{+}|^{2} \cdot [-m(t)] + |\vec{R}_{-}|^{2} \cdot [+m(t)],$$
herized have
$$\mathbf{R}^{-1}(t) = \frac{1}{2} |\vec{R}_{+}|^{2} \cdot [-m(t)] + |\vec{R}_{-}|^{2} \cdot [+m(t)],$$

For unpolarized beam $R = \frac{1}{2} |R_+|^2 + |R_-|^2$ no effect

For ideal mirror, $|R_{+}|=1$, $|R_{-}|=0$, and $H(t) = H_{0}\sin(\omega \cdot t)$ the reflected intensity is modulated as : $R = \frac{\mu_{DW}H_{0}}{\omega \cdot l} \left[-\cos(\omega \cdot t) \right]_{-}^{-}$



Time modulation and smearing of PNR response



$$I(t) = R_0 + \frac{1}{l_s} \int_{-l_s/2}^{l_s/2} dx_s \frac{1}{l_d} \int_{-l_d/2}^{l_d/2} dx_d \int d\lambda R_p \cos\{\omega \cdot [t - (L + x_s + x_d)/v_n]\} \cdot W(\lambda)$$

 $v_n = 1/c\lambda$, $c = (m_n/2\pi\hbar) = 0.2528 \cdot 10^{-5} \text{ sec/(cm} \cdot \text{Å})$, $t = \bar{t} + \Delta t$, $\bar{t} = c\lambda L$, where v_n is neutron velocity, and $W(\lambda)$ is spectral function $(\Delta t/\bar{t}) = \pm (l/L) \pm (\Delta \lambda/\lambda)$ uncertainty in time, If $\lambda = 4 \text{ Å}$, L = 100 cm, l = 1 cm, $\Delta \lambda/\lambda = 1\%$, then $\bar{t} \approx 1 \text{ msec}$,

 $\Delta t_l = c\lambda l \cong 10\mu \,\mathrm{sec}, \quad \Delta t_\lambda \cong 10\mu \,\mathrm{sec}, \quad \Delta t \cong 20\mu \,\mathrm{sec}, \quad f = 50 \,\mathrm{kHz}$

Time modulation with double reflection



$$\begin{split} &I(t) \propto \frac{1}{l} \int_{1/2}^{l/2} dx \int d\lambda \operatorname{R}_{2} [t - (L_{2} + x)/v_{n}] \cdot \operatorname{R}_{1} [t - (L_{1} + L_{2} + x)/v_{n}] \cdot W_{0}(\lambda) \\ &\operatorname{R}_{1} = \operatorname{R}_{0} + \operatorname{R}_{P} \cos \{ \omega_{1} [t - (\tau_{x} + \tau_{2})] \}, \qquad \tau_{x} = x/v_{n}, \quad \tau_{2} = L_{2}/v_{n} \\ &\operatorname{R}_{2} = \operatorname{R}_{0} + \operatorname{R}_{P} \cos \{ \omega_{2} [t - (\tau_{x} + \tau_{2} + \tau_{1})] \}, \qquad \tau_{1} = L_{1}/v_{n} \\ &< \cos\varphi_{1} \cdot \cos\varphi_{2} >= \frac{1}{2} < \cos(\varphi_{1} - \varphi_{2}) + \cos(\varphi_{1} + \varphi_{2}) >\approx \frac{1}{2} < \cos(\varphi_{1} - \varphi_{2}) > \\ &\varphi_{1} - \varphi_{2} = \omega_{H} \cdot (t - \tau_{x}) - [\omega_{H} \cdot \tau_{2} - \omega_{1}\tau_{1}] \approx \omega_{H} \cdot t - \overline{\phi} \\ &\text{if heterodyne frequency } \omega_{H} = \omega_{1} - \omega_{2} \quad \text{is small } \omega_{H}\tau_{x} <<1, \\ &\overline{\phi} = (\omega_{H}L_{2} - \omega_{1}L_{1})/\overline{v_{n}} \text{ and at MIEZE conditions } |\omega_{H}\tau_{2} - \omega_{1}\tau_{1}| << \Delta\lambda/\overline{\lambda} \end{split}$$

Heterodyne setup



PNR data (SuperADAM) + fit



PNR at RESEDA



Hysteresis loop from periodic [FeSi]x40 multilayer



Heterodyne effect measured @ RESEDA (FRM2-Muenchen, 2014)



Heterodyne @ fH=500 Hz for frequences f=75 kHz & 0.4 MHz



Double reflection @ frequences: f1 = 400 kHz, f2=400.5 kHz





Transmission after reflection @ frequences f1 = 400 kHz, f2 400.5 kHz





Heterodyne modulation contrast



Frequency scans at different ac amplitudes



 Δf = + 500 Hz, L₂/L₁ = 20, no bias field

Conclusions

Heterodyne effect is experimentally proven to work with broad wavelength spread

Prototype of the spectrometer fro TRAC PNR is built

Perspectives for development of e.g. spin-echo spectrometer for inelastic PNR are suggested

Frequency dependence of magnetization reversal in thin Fe(100) films

K. Zhernenkov,^{1,*} D. Gorkov,¹ B. P. Toperverg,^{1,2} and H. Zabel¹



Integrated over time PNR from 100 nm Fe film in AC field parallel to DC field and one of easy axes

Frequency dependence of fitting parameters

$$\langle \mathbf{R}^{--} \rangle = \frac{1}{4} \left| R_{+} (1 - \bar{c}_{\gamma}) + R_{-} (1 + \bar{c}_{\gamma}) \right|^{2}$$

 $\langle \mathbf{R}^{++} \rangle = \frac{1}{2} \left| R_{\perp} (1 + \overline{c}_{\perp}) + R_{\perp} (1 - \overline{c}_{\perp}) \right|^2$

$$\langle \mathbf{R}^{+-} \rangle = \langle \mathbf{R}^{-+} \rangle = \frac{1}{4} |\mathbf{R}_{+} - \mathbf{R}_{-}|^{2} \overline{s_{\gamma}^{2}}$$

$$\overline{c}_{\gamma} = \frac{1}{T} \int_{0}^{T} dt \cos \gamma(t) = w_{\uparrow} - w_{\downarrow}$$

$$\overline{s_{\gamma}^2} = \frac{1}{T} \int_0^T dt \sin^2 \gamma(t) = w_{\perp}$$

 $w_{\uparrow} - w_{\downarrow} + w_{\perp} = 1$ $\Delta = \left\langle \cos^2 \gamma \right\rangle - \left\langle \cos \gamma \right\rangle^2 \text{ is dispersion}$



@ 400 kHz crossover from 180°+ 90° DW to 180°DW

General Equation for PNR with 1D polarization analysis

If
$$\vec{P}_{i,f} = \{0, \pm P_{i,f}, 0\}$$
, then $(\vec{P}_{i,f} \cdot \vec{b}) = \pm P_{i,f} \cos \gamma$,
 $R = \frac{1}{4} [R_+ |^2 + |R_-|^2] [+ P_i \cdot P_f \cdot \langle \cos^2 \gamma \rangle]$
 $+ \frac{1}{4} [R_+ |^2 - |R_-|^2] [+ P_i \cdot P_f \cdot \langle \cos \gamma \rangle$
 $+ \frac{1}{2} \operatorname{Re} (R_+^* R_-) P_i \cdot P_f \cdot \langle \sin^2 \gamma \rangle$
In ideal case $P_{i,f} = \pm 1$
Non - spin - flip reflectivities:
 $R^{\pm\pm} = \frac{1}{4} \triangleleft (1 \pm \cos \gamma) \cdot R^+ + (1 \mp \cos \gamma) \cdot R^- |^2 >$
Spin - flip - reflectivities:
 $R^{\pm\mp} = \frac{1}{4} \triangleleft R^+ - R^- |^2 \cdot \sin^2 \gamma >$

where averaging runs over reflecting surface and time

Time averaged PNR for crossed DC and AC fields (with polarization parallel DC field)

